# Calculation of I-adic Local Fourier Transformations 

Lei Fu

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A $\overline{\mathbf{Q}}_{l}$-representation of the galois group of the function field of $X$ unramified on $X$ gives rise to a $\overline{\mathbf{Q}}_{1}$-sheaf on $X$. A complex of such galois representations gives an object in $D\left(X, \overline{\mathbf{Q}}_{l}\right)$.

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Analogue: $\hat{f}\left(t^{\prime}\right)=\int_{-\infty}^{\infty} f(t) e^{i t t^{\prime}} d t$.

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Analogue: $\int_{-\infty}^{\infty} g(x) e^{i f(x) t^{\prime}} d x$, where $g(x)$ is supported in $[a, b]$.

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& =\quad \cdots \\
& =\quad O\left(\frac{1}{t^{\prime n}}\right) \text { for all } n
\end{aligned}
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We would like to have an I-adic analogue of the stationary phase principle.

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\mathcal{F}^{\left(0, \infty^{\prime}\right)}\left([r]_{*}\left(\alpha^{*} \mathcal{L}_{\psi} \otimes \mathcal{K}_{\chi}\right)\right) \cong[r+s]_{*}\left(\beta^{*} \mathcal{L}_{\psi} \otimes \mathcal{K}_{\chi^{-1}} \otimes[s]^{*} \mathcal{K}_{\chi_{2}}\right)
$$

for any multiplicative character $\chi$.
Laumon-Malgrange also made conjectures for the local Fourier transformations $\mathcal{F}^{\left(\infty, 0^{\prime}\right)}$ and $\mathcal{F}^{\left(\infty, \infty^{\prime}\right)}$. Similar results hold for these transformations.

The case $\alpha(\sqrt[r]{t})=\frac{1}{(\sqrt[y]{t})^{s}}$

## The case $\alpha(\sqrt[r]{t})=\frac{1}{(\sqrt[r]{t})^{s}}$

The above system of equations becomes

$$
\left\{\begin{array}{c}
t^{-\frac{s}{r}}+t t^{\prime}=\beta\left(\frac{1}{r+s}\right) \\
-\frac{s}{r} t^{-\frac{s}{r}-1}+t^{\prime}=0
\end{array}\right.
$$

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From the second equation, we get

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t=\left(\frac{r t^{\prime}}{s}\right)^{-\frac{r}{r+s}}
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$$
\beta\left(\frac{1}{\sqrt[r+s]{t^{\prime}}}\right)=\frac{1+\frac{s}{r}}{\left(\frac{s}{r}\right)^{\frac{s}{r+s}}}\left(\sqrt[r+s]{t^{\prime}}\right)^{s}
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Let's prove

$$
[r+s]^{*} \mathcal{F}^{\left(0, \infty^{\prime}\right)}\left([r]_{*} \alpha^{*} \mathcal{L}_{\psi}\right)=[r+s]^{*}[r+s]_{*}\left(\beta^{*} \mathcal{L}_{\psi} \otimes[s]^{*} \mathcal{K}_{\chi_{2}}\right)
$$

Method

Method

$$
\begin{aligned}
& \mathbf{A}^{1}-\{0\} \quad \xrightarrow{\alpha} \mathbf{A}^{1} \\
& \downarrow[r] \\
& \mathbf{A}^{1}-\{0\} \\
& \downarrow j \\
& \mathbf{A}^{1}
\end{aligned}
$$

## Method

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\end{aligned}
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To study $[r+s]^{*} \mathcal{F}^{\left(0, \infty^{\prime}\right)}\left([r]_{*} \alpha^{*} \mathcal{L}_{\psi}\right)$,

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$$

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## Method

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& \downarrow j \\
& \mathbf{A}^{1}
\end{aligned}
$$

To study $[r+s]^{*} \mathcal{F}^{\left(0, \infty^{\prime}\right)}\left([r]_{*} \alpha^{*} \mathcal{L}_{\psi}\right)$, we construct a morphism $f$ such that $[r+s]^{*} \mathcal{F}\left(j_{!}[r]_{*} \alpha^{*} \mathcal{L}_{\psi}\right)$ can be related to $\mathcal{F}\left(R f_{!} \overline{\mathbf{Q}}_{l}\right)$.

## Method

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## Method

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\end{aligned}
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To study $[r+s]^{*} \mathcal{F}^{\left(0, \infty^{\prime}\right)}\left([r]_{*} \alpha^{*} \mathcal{L}_{\psi}\right)$, we construct a morphism $f$ such that $[r+s]^{*} \mathcal{F}\left(j_{!}[r]_{*} \alpha^{*} \mathcal{L}_{\psi}\right)$ can be related to $\mathcal{F}\left(R f_{!} \overline{\mathbf{Q}}_{l}\right)$. At all ramification points, $R f_{!} \overline{\mathbf{Q}}_{l}$ is tame, and the local Fourier transformations for $R f_{!} \overline{\mathbf{Q}}_{/}$can be calculated.

## Study of $[r+s]^{*} \mathcal{F}\left(j![r]_{*} \alpha^{*} \mathcal{L}_{\psi}\right)$

## Study of $[r+s]^{*} \mathcal{F}\left(j_{!}[r]_{*} \alpha^{*} \mathcal{L}_{\psi}\right)$

One can check directly

$$
=\quad \begin{array}{cc}
\operatorname{Tr}\left(F_{t^{\prime}},\left([r+s]^{*} \mathcal{F}\left(j_{!}[r]_{*} \alpha^{*} \mathcal{L}_{\psi}\right)\right)_{\bar{t}^{\prime}}\right) \\
& \sum_{x \in\left(\mathbf{A}^{1}-\{0\}\right)\left(F_{q^{k}}\right)} \psi_{k}\left(x^{r} t^{\prime r+s}+1 / x^{s}\right)
\end{array}
$$

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= & \sum_{x \in\left(\mathbf{A}^{1}-\{0\}\right)\left(\mathbf{F}_{q^{k}}\right)} \psi_{k}\left(t^{\prime s}\left(\left(x t^{\prime}\right)^{r}+1 /\left(x t^{\prime}\right)^{s}\right)\right)
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\end{aligned}
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\end{aligned}
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\begin{array}{cc} 
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\end{array}
$$

that is,

$$
\begin{aligned}
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\end{aligned}
$$

Consider the morphism

$$
f: \mathbf{A}^{1}-\{0\} \rightarrow \mathbf{A}^{1}, f(x)=x^{r}+\frac{1}{x^{s}} .
$$

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$$
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$$
[r+s]^{*} \mathcal{F}\left(j_{!}[r]_{*} \alpha^{*} \mathcal{L}_{\psi}\right) \cong[s]^{*} \mathcal{F}\left(R f_{!} \overline{\mathbf{Q}}_{l}\right)
$$

on $\mathbf{A}^{1}-\{0\}$.

## Study of $\mathcal{F}\left(R f \overline{\mathbf{Q}}_{l}\right)$

Study of $\mathcal{F}\left(R f \overline{\mathbf{Q}}_{l}\right)$
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$$
\left.\left(R f_{!} \overline{\mathbf{Q}}_{l}\right)\right|_{\eta_{y}} \cong\left(\mathbf{Q}_{l} \oplus \mathcal{K}_{\chi_{2}}\right)^{\oplus d} \oplus \overline{\mathbf{Q}}_{l}^{r+s-2 d}
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where $d=(r, s)$.

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$\left.\left(R f_{!} \overline{\mathbf{Q}}_{l}\right)\right|_{\eta_{\infty}}$ is tamely ramified.

By the Stationary Phase Principle, we have
$\left.\left(\mathcal{F}\left(R f_{!} \overline{\mathbf{Q}}_{l}\right)\right)\right|_{\eta_{\infty^{\prime}}}$

By the Stationary Phase Principle, we have

$$
\left.\left.\cong \bigoplus_{y=\frac{1+\frac{s}{r}}{\zeta^{s}\left(\frac{s}{r}\right)^{\frac{s}{r+s}}}} \mathcal{F}^{\left(y, \infty^{\prime}\right)}\left(\left.\left(R f_{!} \overline{\mathbf{Q}}_{l}\right)\right|_{\eta_{y}}\right) \bigoplus \mathcal{F}^{\left(\infty, \infty^{\prime}\right)}\left(\left(R f_{!} \overline{\mathbf{Q}}_{l}\right)\right)\right|_{\left.\left.\eta_{\infty^{\prime}}\right)_{\eta_{\infty}}\right)}\right)
$$

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$$
\begin{aligned}
& \left.\left(\mathcal{F}\left(R f_{!} \overline{\mathbf{Q}}_{l}\right)\right)\right|_{\eta_{\infty^{\prime}}} \\
\cong & \bigoplus_{y=\frac{1+\frac{s}{r}}{\zeta^{s}\left(\frac{s}{r}\right)^{\frac{s}{r}+s}}} \mathcal{F}^{\left(y, \infty^{\prime}\right)}\left(\left.\left(R f_{!} \overline{\mathbf{Q}}_{l}\right)\right|_{\eta_{y}}\right) \bigoplus \mathcal{F}^{\left(\infty, \infty^{\prime}\right)}\left(\left(R f_{!} \overline{\mathbf{Q}}_{l}\right)_{\eta_{\infty}}\right) \\
\cong & \left.\bigoplus_{\zeta^{s}} \mathcal{F}^{\left(\frac{1+\frac{s}{r}}{\zeta^{s}\left(\frac{s}{r}\right)^{\frac{s}{r+s}}}, \infty^{\prime}\right.}\right)\left(\mathcal{K}_{\chi_{2}}^{\oplus d}\right)
\end{aligned}
$$

By the Stationary Phase Principle, we have

$$
\begin{aligned}
&\left.\left(\mathcal{F}\left(R f_{!} \overline{\mathbf{Q}}_{l}\right)\right)\right|_{\eta_{\infty^{\prime}}} \\
& \cong \bigoplus_{y=\frac{1+\frac{s}{r}}{\zeta^{s}\left(\frac{s}{r} \frac{s}{r+s}\right.}} \mathcal{F}^{\left(y, \infty^{\prime}\right)}\left(\left.\left(R f_{!} \overline{\mathbf{Q}}_{l}\right)\right|_{\eta_{y}}\right) \bigoplus \mathcal{F}^{\left(\infty, \infty^{\prime}\right)}\left(\left(R f_{!} \overline{\mathbf{Q}}_{l}\right)_{\eta_{\infty}}\right) \\
& \cong\left.\bigoplus_{\zeta^{s}} \mathcal{F}^{\left(\frac{1+\frac{s}{r}}{\zeta^{s}\left(\frac{s}{r}\right)^{\frac{s}{r+s}}}, \infty^{\prime}\right.}\right) \\
&\left(\mathcal{K}_{\chi_{2}}^{\oplus d}\right) \\
& \cong \bigoplus_{\zeta^{s}}\left(\mathcal{L}_{\psi}\left(\frac{1+\frac{s}{r}}{\zeta^{s}\left(\frac{s}{r}\right)^{\frac{s}{r+s}}} t^{\prime}\right) \otimes \mathcal{K}_{\chi_{2}}\right)^{\oplus d},
\end{aligned}
$$

that is,

$$
\left(\mathcal{F}\left(R f \overline{\mathbf{Q}}_{l}\right)\right)_{\eta_{\infty^{\prime}}} \cong \bigoplus_{\zeta^{s}}\left(\mathcal{L}_{\psi}\left(\frac{1+\frac{s}{r}}{\zeta^{s}\left(\frac{s}{r}\right)^{\frac{s}{r+s}}} t^{\prime}\right) \otimes \mathcal{K}_{\chi_{2}}\right)^{\oplus d} .
$$

By the Stationary Phase Principle again, we have
that is,

$$
\left(\mathcal{F}\left(R f \overline{\mathbf{Q}}_{l}\right)\right)_{\eta_{\infty^{\prime}}} \cong \bigoplus_{\zeta^{s}}\left(\mathcal{L}_{\psi}\left(\frac{1+\frac{s}{r}}{\zeta^{s}\left(\frac{s}{r}\right)^{\frac{s}{r+s}}} t^{\prime}\right) \otimes \mathcal{K}_{\chi_{2}}\right)^{\oplus d} .
$$

By the Stationary Phase Principle again, we have

$$
\left.[r+s]^{*} \mathcal{F}^{\left(0, \infty^{\prime}\right)}\left([r]_{*} \alpha^{*} \mathcal{L}_{\psi}\right) \cong\left([r+s]^{*} \mathcal{F}\left(j![r]_{*} \alpha^{*} \mathcal{L}_{\psi}\right)\right)\right|_{\eta_{\infty^{\prime}}}
$$

that is,

$$
\left(\mathcal{F}\left(R f \overline{\mathbf{Q}}_{l}\right)\right)_{\eta_{\infty^{\prime}}} \cong \bigoplus_{\zeta^{s}}\left(\mathcal{L}_{\psi}\left(\frac{1+\frac{s}{r}}{\zeta^{s}\left(\frac{s}{r}\right)^{\frac{s}{r+s}}} t^{\prime}\right) \otimes \mathcal{K}_{\chi_{2}}\right)^{\oplus d} .
$$

By the Stationary Phase Principle again, we have

$$
\begin{aligned}
{[r+s]^{*} \mathcal{F}^{\left(0, \infty^{\prime}\right)}\left([r]_{*} \alpha^{*} \mathcal{L}_{\psi}\right) } & \left.\cong\left([r+s]^{*} \mathcal{F}\left(j_{j}[r]_{*} \alpha^{*} \mathcal{L}_{\psi}\right)\right)\right|_{\eta_{\infty^{\prime}}} \\
& \left.\cong\left([s]^{*} \mathcal{F}\left(R f \overline{\mathbf{Q}}_{l}\right)\right)\right|_{\eta_{\infty^{\prime}}}
\end{aligned}
$$

that is,

$$
\left(\mathcal{F}\left(R f_{!} \overline{\mathbf{Q}}_{l}\right)\right)_{\eta_{\infty^{\prime}}} \cong \bigoplus_{\zeta^{s}}\left(\mathcal{L}_{\psi}\left(\frac{1+\frac{s}{r}}{\zeta^{s}\left(\frac{s}{r}\right)^{\frac{s}{r+s}}} t^{\prime}\right) \otimes \mathcal{K}_{\chi_{2}}\right)^{\oplus d}
$$

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$$
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& \left.\cong\left([s]^{*} \mathcal{F}\left(R f_{!} \overline{\mathbf{Q}_{l}}\right)\right)\right|_{\eta_{\infty^{\prime}}} \\
& \cong \bigoplus_{\zeta^{s}}[s]^{*}\left(\mathcal{L}_{\psi}\left(\frac{1+\frac{s}{r}}{\zeta^{s}\left(\frac{s}{r}\right)^{\frac{s}{r+s}}} t^{\prime}\right) \otimes \mathcal{K}_{\chi_{2}}\right)^{\oplus d}
\end{aligned}
$$

that is,

$$
\left(\mathcal{F}\left(R f_{!} \overline{\mathbf{Q}}_{l}\right)\right)_{\eta_{\infty^{\prime}}} \cong \bigoplus_{\zeta^{s}}\left(\mathcal{L}_{\psi}\left(\frac{1+\frac{s}{r}}{\zeta^{s}\left(\frac{s}{r}\right)^{\frac{s}{r+s}}} t^{\prime}\right) \otimes \mathcal{K}_{\chi_{2}}\right)^{\oplus d}
$$

By the Stationary Phase Principle again, we have

$$
\begin{aligned}
{[r+s]^{*} \mathcal{F}^{\left(0, \infty^{\prime}\right)}\left([r]_{*} \alpha^{*} \mathcal{L}_{\psi}\right) } & \cong\left(\left.[r+s]^{*} \mathcal{F}\left(j\left[[r] * \alpha^{*} \mathcal{L}_{\psi}\right)\right)\right|_{\eta_{\infty^{\prime}}}\right. \\
& \left.\cong\left([s]^{*} \mathcal{F}\left(R f \overline{\mathbf{Q}}_{l}\right)\right)\right|_{\eta_{\infty^{\prime}}} \\
& \cong \bigoplus_{\zeta^{s}}[s]^{*}\left(\mathcal{L}_{\psi}\left(\frac{1+\frac{s}{r}}{\zeta^{s}\left(\frac{s}{r}\right)^{\frac{s}{r}+s}} t^{\prime}\right) \otimes \mathcal{K}_{\chi_{2}}\right)^{\oplus d} \\
& \cong \bigoplus_{\zeta^{s}}\left(\mathcal{L}_{\psi}\left(\frac{1+\frac{s}{r}}{\zeta^{s}\left(\frac{s}{r} \frac{s}{r}\right)^{\frac{s}{r+s}} t^{\prime s}}\right) \otimes[s]^{*} \mathcal{K}_{\chi_{2}}\right)^{\oplus d}
\end{aligned}
$$

that is,

$$
\left(\mathcal{F}\left(R f_{!} \overline{\mathbf{Q}}_{l}\right)\right)_{\eta_{\infty^{\prime}}} \cong \bigoplus_{\zeta^{s}}\left(\mathcal{L}_{\psi}\left(\frac{1+\frac{s}{r}}{\zeta^{s}\left(\frac{s}{r}\right)^{\frac{s}{r+s}}} t^{\prime}\right) \otimes \mathcal{K}_{\chi_{2}}\right)^{\oplus d}
$$

By the Stationary Phase Principle again, we have

$$
\begin{aligned}
{[r+s]^{*} \mathcal{F}^{\left(0, \infty^{\prime}\right)}\left([r]_{*} \alpha^{*} \mathcal{L}_{\psi}\right) } & \cong\left(\left.[r+s]^{*} \mathcal{F}\left(j\left[[r]_{*} \alpha^{*} \mathcal{L}_{\psi}\right)\right)\right|_{\eta_{\infty^{\prime}}}\right. \\
& \left.\cong\left([s]^{*} \mathcal{F}\left(R f \overline{\mathbf{Q}}_{1}\right)\right)\right|_{\eta_{\infty^{\prime}}} \\
& \cong \bigoplus_{\zeta^{s}}[s]^{*}\left(\mathcal{L}_{\psi}\left(\frac{1+\frac{s}{r}}{\zeta^{s}\left(\frac{s}{r}\right)} t^{\frac{s}{r}+s} t^{\prime}\right) \otimes \mathcal{K}_{\chi_{2}}\right)^{\oplus d} \\
& \cong \bigoplus_{\zeta^{s}}\left(\mathcal{L}_{\psi}\left(\frac{1+\frac{s}{r}}{\zeta^{s}\left(\frac{s}{r}\right.}{ }^{\frac{s}{r}+5} t^{\prime s}\right) \otimes[s]^{*} \mathcal{K}_{\chi_{2}}\right)^{\oplus d} \\
& \cong[r+s]^{*}[r+s]_{*}\left(\beta^{*} \mathcal{L}_{\psi} \otimes[s]^{*} \mathcal{K}_{\chi_{2}}\right) .
\end{aligned}
$$

The proof of the Laumon-Malgrange conjecture for general $\alpha$ involves more complicated changes of variables,

The proof of the Laumon-Malgrange conjecture for general $\alpha$ involves more complicated changes of variables, and a subtle use of a fact stronger than the Stationary Phase Principle.

